



Development of **MH4D** for Simulation of Emerging Concept Fusion Experiments

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Main Ideas



- PSI-Center: further developing a tetrahedral mesh MHD simulation code to model Emerging Concept (EC) experiments (**MH4D**: Lionello & Schnack, 2004)
 - represent the complex 3-dimensional geometry of many EC experiments
 - using existing programs for gridding, parallel implementation, and visualization
- Current status
 - vacuum fields in code
 - input/output in General Mesh Viewer (GMV) format (Ortega, 2000)
 - preliminary simulations of HIT-SI and ZAP experiments with near experimental parameters

Code: Resistive MHD Model



$$\frac{\partial \mathbf{A}}{\partial t} = -\mathbf{E} \quad \mathbf{B} = \nabla \times \mathbf{A} \quad \mathbf{J} = \nabla \times \mathbf{B}$$

$$\mathbf{E} = -\mathbf{v} \times \mathbf{B} + \eta \mathbf{J}$$

$$\frac{\partial \rho}{\partial t} = -\nabla \cdot \rho \mathbf{v}$$

$$\rho \left(\frac{\partial \mathbf{v}}{\partial t} + \mathbf{v} \cdot \nabla \mathbf{v} \right) = -\nabla p + \mathbf{J} \times \mathbf{B} + \nu \nabla^2 \mathbf{v}$$

$$\frac{\partial p}{\partial t} + \mathbf{v} \cdot \nabla p = -\gamma p \nabla \cdot \mathbf{v} + \nabla \cdot \mathbf{q} + Q$$

Alfvén velocity: $V_A = B/\sqrt{\rho}$

Alfvén transit time: $\tau_A = (2\pi r_{axis})/V_A$

Resistive diffusion time : $\tau_R = 1/(\eta\lambda^2)$

Lundquist number: $S = \tau_R/\tau_A$

Include Vacuum Fields



Assume external fields:

$$\dot{\mathbf{A}} = -\mathbf{E}_{int}$$

$$\text{and } \mathbf{B}_{total} = \mathbf{B} + \mathbf{B}_{ext}$$

$$\Rightarrow \dot{\mathbf{A}} = \mathbf{v} \times (\mathbf{B} + \mathbf{B}_{ext}) - \eta \mathbf{J} + \mathbf{E}_{ext}$$

$$\Rightarrow \rho \left(\frac{D\mathbf{v}}{Dt} \right) = -\nabla p + \mathbf{J} \times (\mathbf{B} + \mathbf{B}_{ext}) + \nu \nabla^2 \mathbf{v}$$

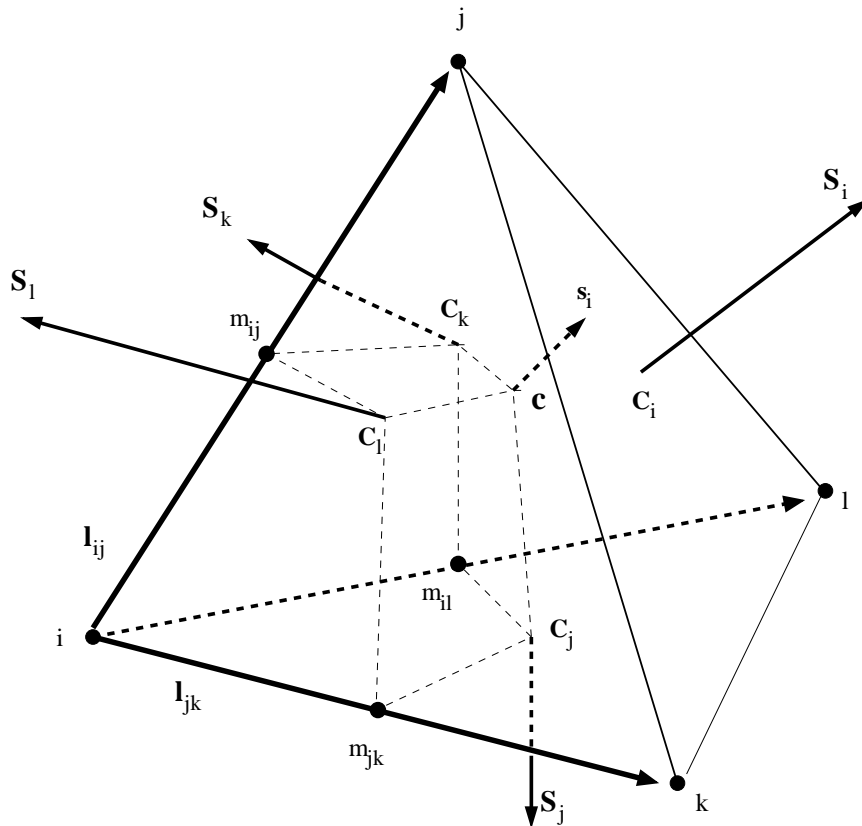
where

$$\mathbf{E}_{ext} = \text{Volt}_{coeff} (\mathbf{v}_{f1} \sin(\omega t) + \mathbf{v}_{f2} \cos(\omega t))$$

$$\mathbf{B}_{ext} = \text{Flux}_{coeff} (\mathbf{v}_{f1} \sin(\omega t) + \mathbf{v}_{f2} \cos(\omega t))$$

If \mathbf{B}_{ext} is time dependent then its induced electric field must be included in \mathbf{E}_{ext} .

Tetrahedral Grid



- Sides are labeled by index of their opposite vertex
- c : centroid of tetrahedron
- m_{ij} : midpoint of edge l_{ij}
- C_i : centroid of side i
- S_i : vector area of side i
- s_i : vector area of dual median surface
- V_c : volume of tetrahedron c

$$\mathbf{s}_i = \frac{1}{3} \mathbf{S}_i$$

$$V_c = \frac{1}{6} \mathbf{l}_{ij} \cdot (\mathbf{l}_{ik} \times \mathbf{l}_{il}) = -\frac{1}{3} \mathbf{l}_{ij} \cdot \mathbf{S}_j$$

Finite Volume Formalism



Integral relations to define differential operators

Gradient:

$$\int_V \nabla f dV = \int_S \hat{\mathbf{n}} f dS$$

Divergence:

$$\int_V \nabla \cdot \mathbf{F} dV = \int_S \hat{\mathbf{n}} \cdot \mathbf{F} dS$$

Curl in 2D

$$\int_s \mathbf{n} \cdot \nabla \times \mathbf{F} dS = \oint_C \mathbf{F} \cdot d\mathbf{l}$$

Curl in 3D

$$\int_V \nabla \times \mathbf{F} dV = \int_S \hat{\mathbf{n}} \times \mathbf{F} dS$$

Magnetic Field



Vector potential \mathbf{A} defined on vertices

- \mathbf{A} varies linearly within tetrahedron

$$\int_V \mathbf{B} dV = \int_V \nabla \times \mathbf{A} dV = \int_S d\mathbf{S} \times \mathbf{A}$$

$$\Rightarrow B_\tau V_\tau = \sum_{s=1}^4 \mathbf{S}_s \times \bar{\mathbf{A}}_s$$

In terms of vertices:

$$\mathbf{B}_\tau = -\frac{1}{3V_\tau} \sum_{\nu(\tau)} \mathbf{S}_\nu \times \mathbf{A}_\nu$$

\mathbf{B}_τ : constant within tetrahedron τ

$\bar{\mathbf{A}}_s$: average of \mathbf{A} over the 3 vertices of side S

V_τ : volume of tetrahedron τ

$\nu(\tau)$: indices of the 4 vertices of tetrahedron τ

\mathbf{S}_ν : outward directed side opposite vertex ν

Divergence of \mathbf{B}



Using Gauss' theorem to dual median volume element surrounding vertex ν

$$\begin{aligned}(\nabla \cdot \mathbf{B})_{\nu} v_{\nu} &= \frac{1}{3} \sum_{\tau} \mathbf{B}_{\tau} \cdot \mathbf{S}_{\nu(\tau)} \\ &= \frac{1}{3} \sum_{\tau} \frac{1}{V_{\tau}} \sum_{\gamma(\tau)} \bar{\mathbf{A}}_{\gamma} \cdot (\mathbf{S}_{\nu} \times \mathbf{S}_{\gamma})\end{aligned}$$

$v_{\nu} = \sum_{\tau(\nu)} \frac{1}{3} V_{\tau}$: volume of dual median volume element

$\tau(\nu)$: indices of the tetrahedra sharing vertex ν

$\nu(\tau)$: index of the side of tetrahedron τ opposite vertex ν

$\gamma(\tau)$: indices of the 4 sides of tetrahedron τ

- Contributions from common sides of adjoining tetrahedra cancel

$$(\nabla \cdot \mathbf{B})_{\nu} = 0$$

Another Derivation of \mathbf{B}



\mathbf{A} varies linearly within tetrahedron

$$\mathbf{A}(\mathbf{x}) = \sum_{\nu} \mathbf{A}_{\nu} \left(1 - \frac{1}{3V_{\tau}} \mathbf{S}_{\nu} \cdot (\mathbf{x} - \mathbf{x}_{\nu}) \right)$$

Take curl of \mathbf{A}

$$B_{\alpha} = \varepsilon_{\alpha\beta\gamma} \frac{\partial A_{\gamma}}{\partial x_{\beta}}$$
$$\Rightarrow \mathbf{B}_{\tau} = -\frac{1}{3V_{\tau}} \sum_{\nu(\tau)} \mathbf{S}_{\nu} \times \mathbf{A}_{\nu}$$

Same as finite volume representation.

Current Density



Apply Stoke's to dual median volume element

$$\int_V \mathbf{J} dV = \int_V \nabla \times \mathbf{B} dV = \int_S \hat{\mathbf{n}} \times \mathbf{B} dS$$

Use \mathbf{B} in terms of \mathbf{A}

$$\mathbf{J}_\nu v_\nu = \frac{1}{3} \sum_{\tau(\nu)} \mathbf{S}_{\nu(\tau)} \times \mathbf{B}_\tau$$

$$\mathbf{J}_\nu v_\nu = \sum_{\tau(\nu)} \sum_{\nu'(\tau)} \mathfrak{M}^\tau(\nu, \nu') \cdot \mathbf{A}_{\nu'}$$

$$\mathfrak{M}^\tau(\nu, \nu') = \frac{1}{9V_\tau} [(\mathbf{S}_{\nu(\tau)} \cdot \mathbf{S}_{\nu'(\tau)}) \mathbf{I} - \mathbf{S}_{\nu'(\tau)} \mathbf{S}_{\nu(\tau)}^T]$$

\mathfrak{M}^τ is symmetric, self adjoint and the discrete “curl-curl” operator

Variational Principle (VP)



- begin with functional: $I(\mathbf{A}) = \frac{1}{2} [(\nabla \times \mathbf{A})^2 - 2\mathbf{J} \cdot \mathbf{A}] dV$
- minimize I w.r.t variations $\delta\mathbf{A}$, let $\mathbf{A} \rightarrow \mathbf{A} + \epsilon\delta\mathbf{A}$

$$I(\mathbf{A} + \epsilon\delta\mathbf{A}) = \frac{1}{2} \int [(\nabla \times \mathbf{A})^2 - 2\mathbf{J} \cdot \mathbf{A}] dV$$
$$+ \epsilon \underbrace{\int (\nabla \times \mathbf{A} \cdot \nabla \times \delta\mathbf{A} - \mathbf{J} \cdot \delta\mathbf{A}) dV}_C$$
$$+ \frac{\epsilon^2}{2} \int (\nabla \times \delta\mathbf{A})^2 dV$$

I has minimum at $\epsilon = 0$ only if $C = 0$. Vector identities and integration by parts imply

$$\int \delta\mathbf{A} \cdot (\nabla \times \nabla \times \mathbf{A} - \mathbf{J}) dV = - \int_S \delta\mathbf{A}_t \cdot (\mathbf{B} \times \hat{\mathbf{n}}) dS$$

$$\text{where } \mathbf{A}_t = (\mathbf{I} - \hat{\mathbf{n}}\hat{\mathbf{n}}) \cdot \mathbf{A}$$

- Let: $I(\mathbf{A}) = \frac{1}{2} [(\nabla \times \mathbf{A})^2 - 2\mathbf{J} \cdot \mathbf{A}] dV$
- Minimum of $I(\mathbf{A})$ is solution to $\nabla \times \nabla \times \mathbf{A} = \mathbf{J}$ with \mathbf{A}_t specified on boundary
- Minimize $I(\mathbf{A})$ on tetrahedral grid with N_ν vertices and N_τ tetrahedra.
 - Expand $\mathbf{A}(\mathbf{x})$ with basis $\alpha_\nu(\mathbf{x})$: $\mathbf{A}(\mathbf{x}) = \sum \mathbf{A}_\nu \alpha_\nu(\mathbf{x})$
 - Expand current density with delta functions:

$$\mathbf{J}(\mathbf{x}) = \sum_\nu v_\nu J_\nu \delta(\mathbf{x} - \mathbf{x}_\nu)$$

VP gives discrete result



Discrete functional:

$$I = \frac{1}{2} \sum_{\nu} \sum_{\nu'} \mathbf{A}_{\nu} \cdot \mathfrak{M}(\nu, \nu') \cdot \mathbf{A}_{\nu'} - \sum_{\nu} \sum_{\nu'} v_{\nu} \mathbf{J}_{\nu} \cdot \mathbf{A}_{\nu'} \alpha_{\nu'}(\mathbf{x}_{\nu})$$

$$\text{where } \mathfrak{M} = \sum_{\tau} \mathfrak{M}^{\tau}$$

To minimize, $\frac{\partial I}{\partial \mathbf{A}_{\mu}} = 0, \quad \forall \mu = 1, \dots, N_{\nu}$

$$\iff \sum_{\nu} \alpha_{\mu}(\mathbf{x}_{\nu}) v_{\nu} \mathbf{J}_{\nu} = \sum_{\tau} \sum_{\nu} \mathfrak{M}^{\tau}(\mu, \nu) \cdot \mathbf{A}_{\nu}, \quad \forall \mu = 1, \dots, N_{\nu}$$

Setting $\alpha_{\mu}(\mathbf{x}_{\nu}) = \delta_{\mu\nu}$ is the finite volume expression for solving \mathbf{J} from \mathbf{A}

Boundary Condition



- Natural boundary condition,

$$\delta \mathbf{A}_t \cdot (\mathbf{B} \times \hat{\mathbf{n}}) = 0$$

- We must specify $\mathbf{A}_t = (\mathbf{I} - \hat{\mathbf{n}}\hat{\mathbf{n}}) \cdot \mathbf{A}$ on the boundary (Dirichlet)
- Item for further PSI-center development

Partial Implicit Time Stepping



Large spectrum of eigenvalues associated with MHD operators
⇒ large range of time scales

$$\frac{\partial u}{\partial t} = \underbrace{\mathcal{M}\{u\}}_{\text{Full MHD Operator}} = \underbrace{\mathcal{F}\{u\}}_{\substack{\text{Fast time scales} \\ \text{Alfvén waves, soundwaves}}} + \underbrace{\mathcal{S}\{u\}}_{\substack{\text{Slow time scales} \\ \text{interesting physics}}}$$

To avoid time step restriction, treat "fast" part of operator implicitly

$$\frac{u^{n+1} - u^n}{\Delta t} = F u^{n+1} + S u^n$$

Quite difficult to achieve precise discrete representation

No need for operator S



F and M are known operators in MHD computations

$$\begin{aligned}\frac{u^{n+1} - u^n}{\Delta t} &= Fu^{n+1} + \underbrace{(M - F)}_{S=M-F} u^n \\ &= Mu^n + \Delta t F \left(\frac{u^{n+1} - u^n}{\Delta t} \right)\end{aligned}$$

Works for arbitrary operator F

Semi-Implicit Method



For an arbitrary operator G

$$\left(I - \underbrace{\Delta t G}_{\text{SI operator}} \right) u^{n+1} = \underbrace{(I + \Delta t M) u^n}_{\text{Explicit}} - \underbrace{\Delta t G}_{\text{SI operator}} u^n$$

Want to choose G for ease of inversion and inclusion of modes of interest

Semi-implicit MHD operator



- use of linearized, ideal MHD wave equation

$$\rho_0 \frac{\partial^2 \mathbf{v}}{\partial t^2} = \nabla \times \nabla \times (\mathbf{v} \times \mathbf{B}_0) \times \mathbf{B}_0$$

- large spectrum of normal modes
- anisotropic spatial operator
- modify equation of motion

$$\rho \frac{D\mathbf{v}}{Dt} = -\nabla p + \mathbf{J} \times \mathbf{B} + \alpha \Delta t \mathbf{S} \left(\frac{\partial \mathbf{v}}{\partial t} \right)$$

$$\text{where } \mathbf{S}(\mathbf{v}) = \nabla \times \nabla \times (\mathbf{v} \times \mathbf{B}) \times \mathbf{B}$$

S is self-adjoint.

Discrete Semi-Implicit Operator



Use variational principle:

$$I(\mathbf{v}) = \frac{1}{2} \int [|\nabla \times (\mathbf{v} \times \mathbf{B})|^2 + 2\mathbf{S} \cdot \mathbf{v}] dV$$

Let:

$$\mathbf{v} \times \mathbf{B} = \sum_j (\mathbf{v}_j \times \mathbf{B}_j) \alpha_j(\mathbf{x}), \quad \mathbf{v} = \sum_j \mathbf{v}_j \alpha_j(\mathbf{x}), \quad \mathbf{S} = \sum_j \mathbf{S}_j \delta(\mathbf{x} - \mathbf{x}_j)$$

Discrete variational principle:

$$I(\mathbf{v}) = \frac{1}{2} \sum_k \sum_j \mathbf{v}_j \cdot \mathfrak{N}(j, k) \cdot \mathbf{v}_k + \mathbf{S}_k \cdot \mathbf{v}_k$$

where $\mathfrak{N}(j, k) = \mathbf{B}_j \times \mathfrak{M}(j, k) \times \mathbf{B}_k$ and is symmetric

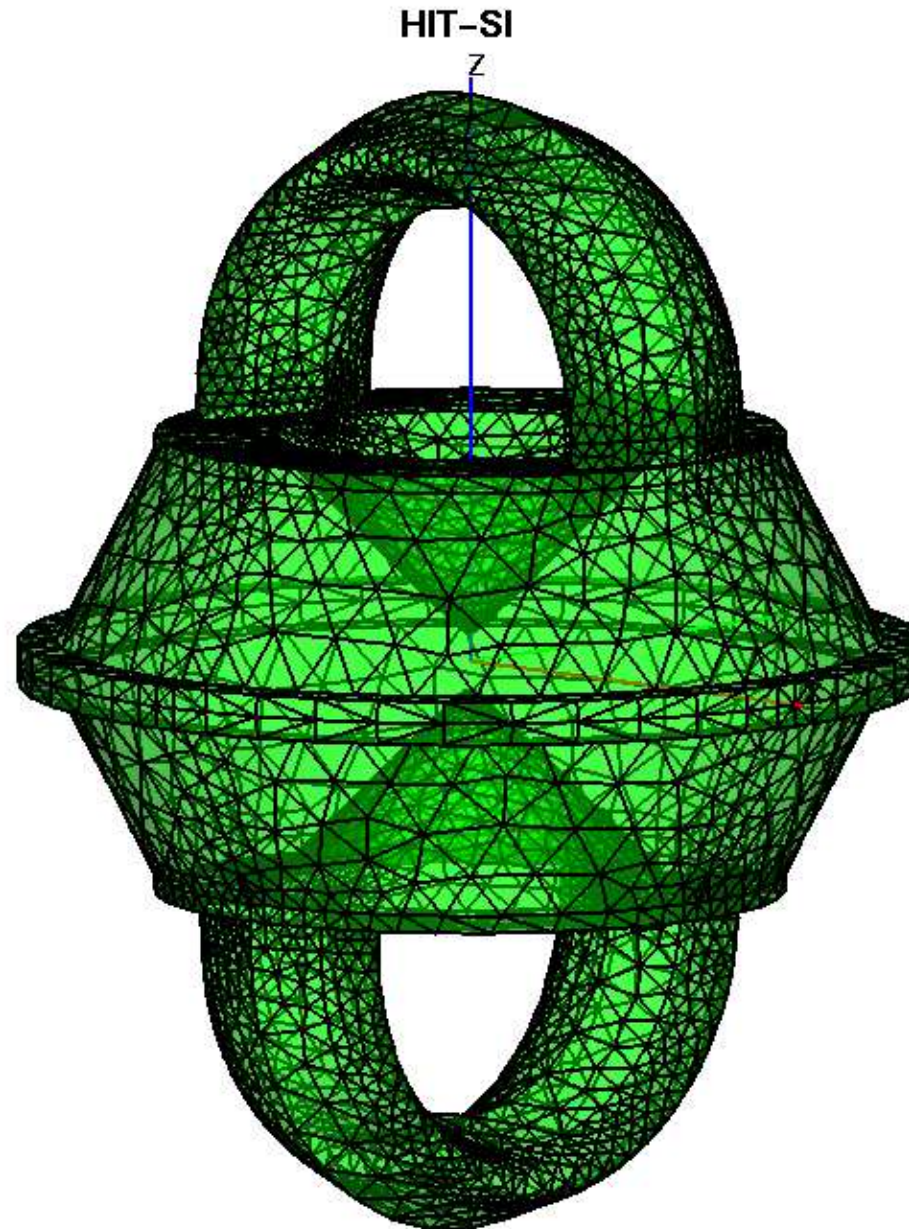
Minimization of $I \rightarrow$ discrete semi-implicit operator \mathbf{S}_i :

$$\frac{\partial I}{\partial \mathbf{v}_i} = 0 \iff \mathbf{S}_i = - \sum_j \mathfrak{N}(i, j) \cdot \mathbf{v}_j, \quad \forall i.$$

Tetrahedral Mesh for HIT-SI

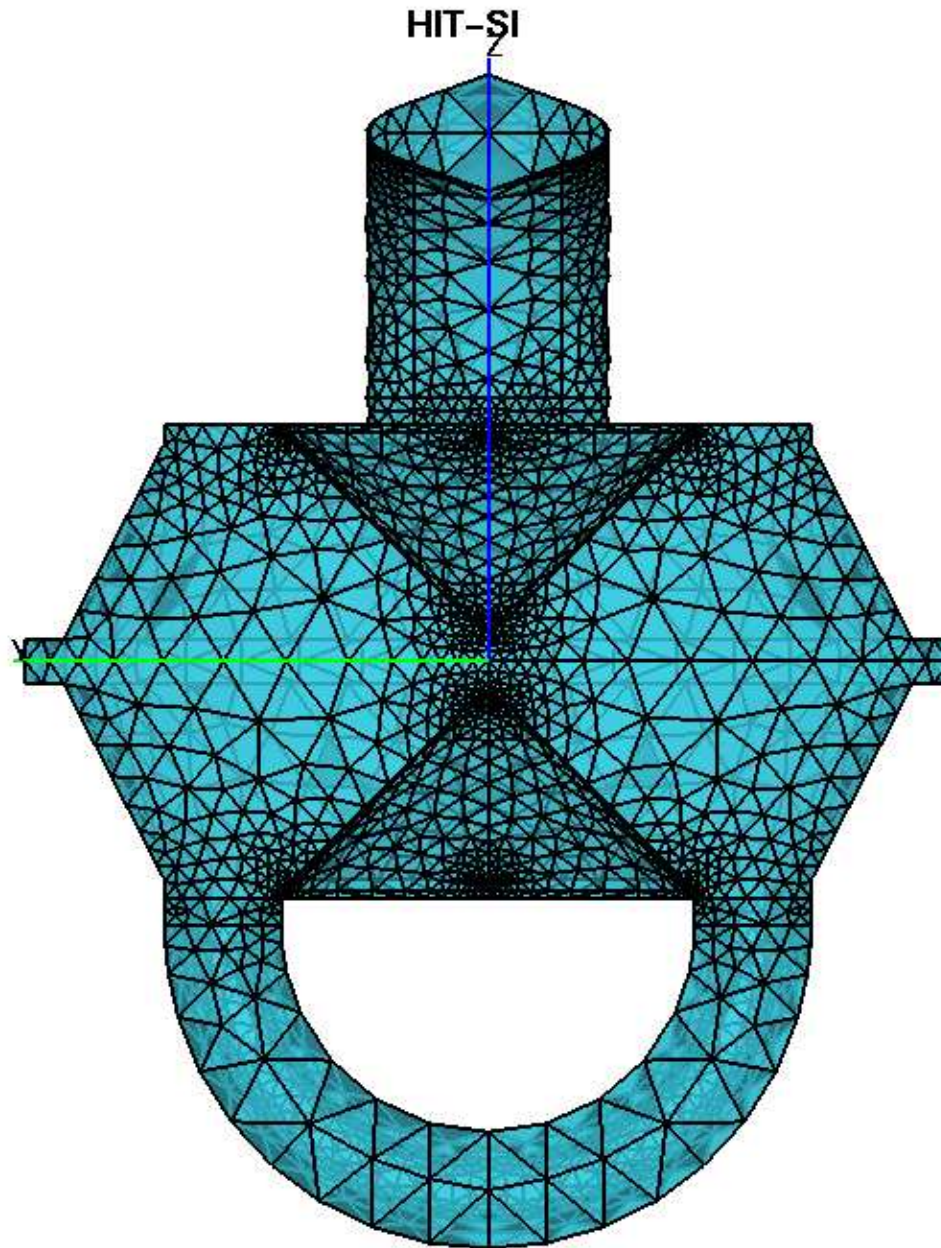


0



HIT-SI: y-z cut plane

1

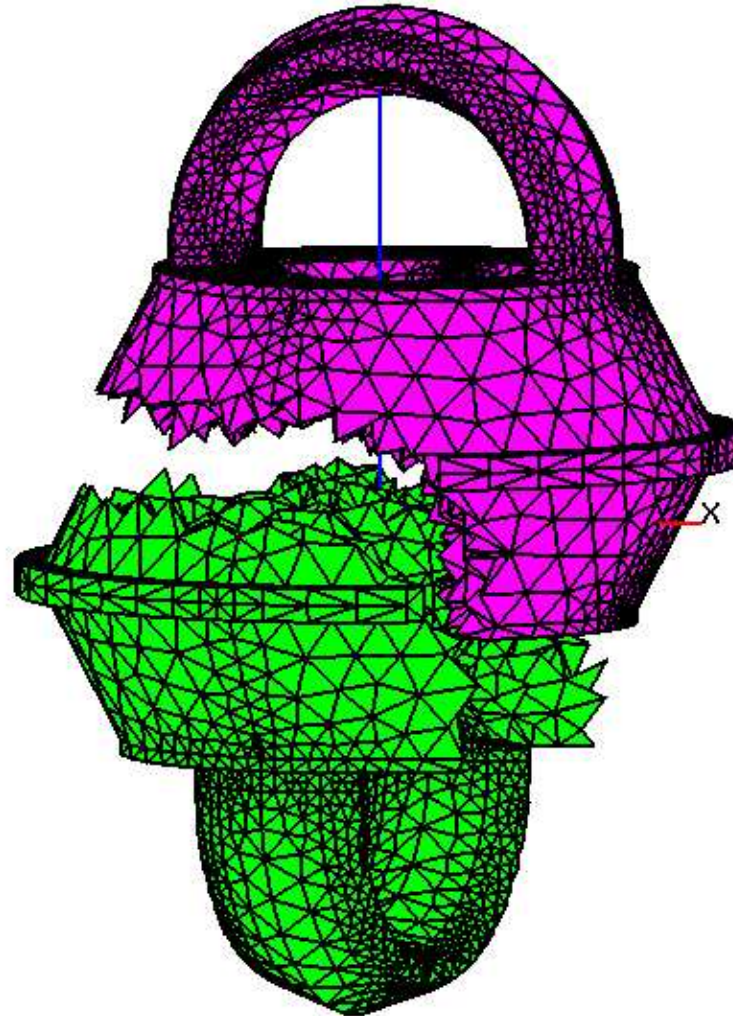


Interior

Domain decomposition

0

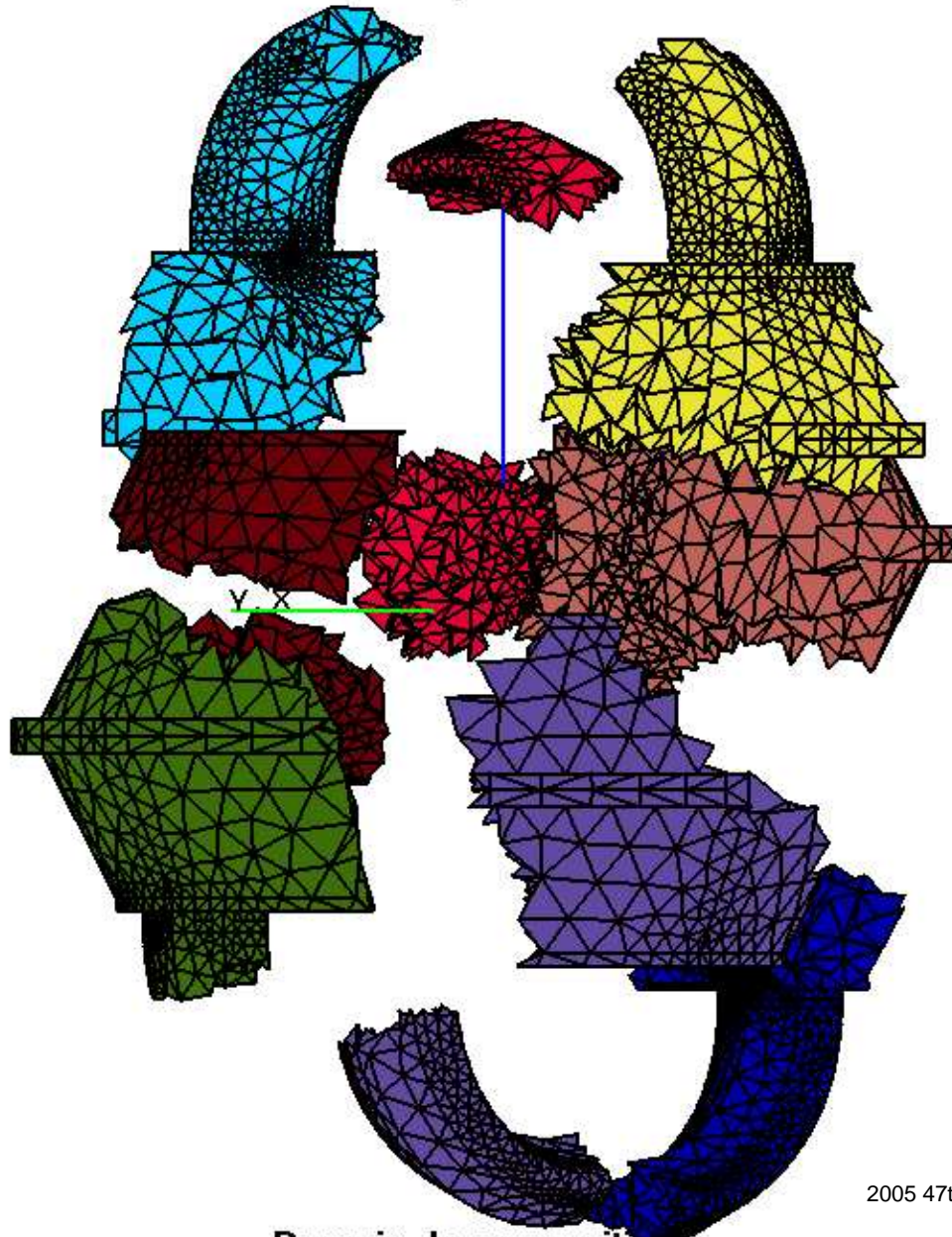
HIT-SI mesh on 2 processors



Domain decomposition

0

HIT-SI on 8 processors

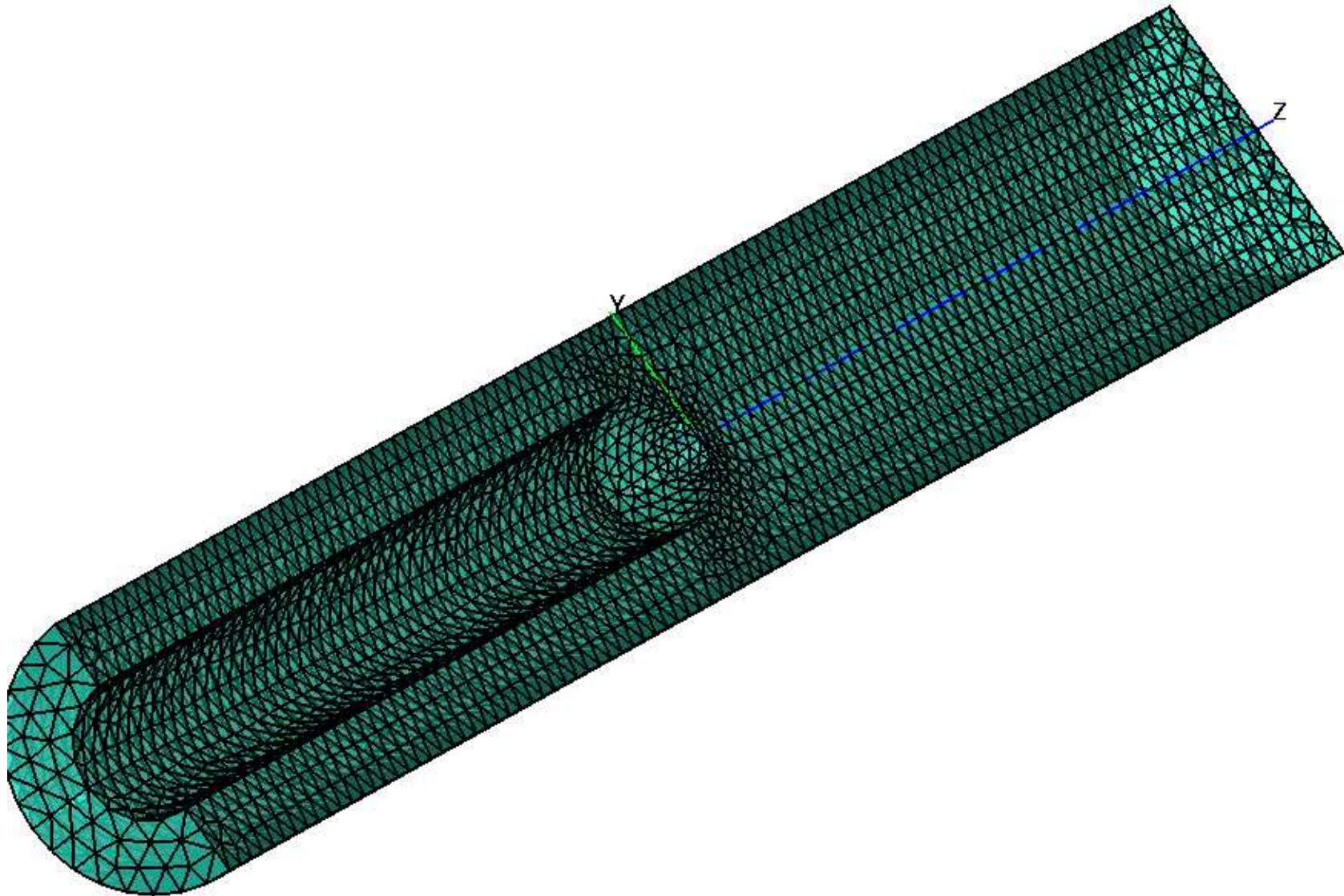


Tetrahedral mesh (1cm)



0

ZaP

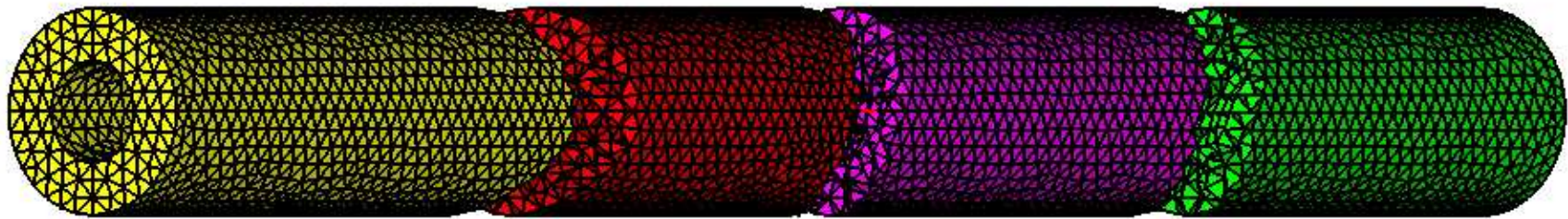


Domain decomposition



0

ZaP on 4 processors



Simulation



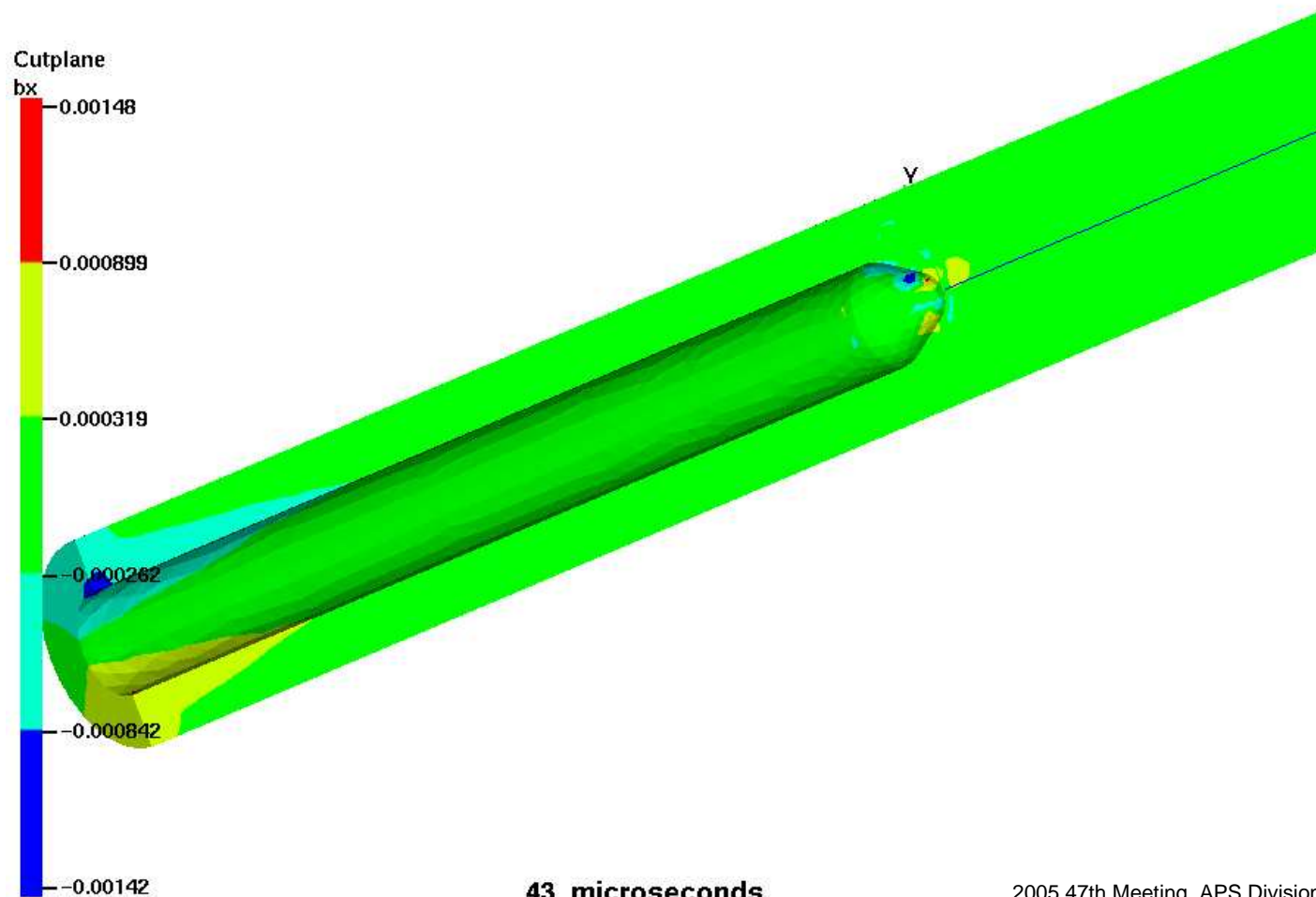
$$T = 10eV, n = 10^{-15} / cm^{-3}$$

Applied 2kV from outer to inner electrode

225285

Formation of B field

4.300011e-05



MH4D computing



Tetrahedral mesh generation:

- T3D (Ryp1, 2004)

Parallel implementations:

- MPI
- METIS (Karypis & Kumar, 1999) for domain decomposition
- PETSc (Baley et al., 2000) for preconditioned CG solver on unstructured grid

Visualization:

- General Mesh Viewer (GMV) (Ortega, 2000) for visualizing data on tetrahedral grid

Conclusion



The PSI-Center is refining the tetrahedral mesh MHD simulation code, **MH4D**, to model Innovative Confinement Concepts experiments.

Current refinement:

- GMV format
- vacuum fields

Current challenges:

- complicated boundary conditions
- neutral gas injection and ionization

Future challenges:

- regions of near zero density